Discrete Mechanics and Optimal Control: Structure preserving integration for the optimal control of mechanical systems

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A Optimal control problem

 $\min_{x,u} J(x,u)$ s.t. $\dot{x}(t) = f(x(t), u(t)), x(0) = x^0$

 well established numerical methods <u>indirect:</u> solve necessary optimality conditions <u>direct:</u> discretize problem *in a clever way* and solve optimization problem



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- well established numerical methods <u>indirect:</u> solve necessary optimality conditions <u>direct:</u> discretize problem *in a clever way* and solve optimization problem
- B **Structure preserving integration** approximate solution x(t) of $\dot{x}(t) = f(x(t))$ via a discrete solution $\{x_k\}_{k=0}^N$ with e.g. $x_{k+1} = \Psi_h(x_k)$
- preserve conservation properties in discrete solution
- geometric integrators, symplectic integrators, here: variational integrators



A+B optimal control problem + structure preserving integration



A+B optimal control problem + structure preserving integration

?

- What is preservation in presence of control u(t)?
- Which properties can be derived from integration theory?
- What do we gain for optimal control problems?



Outline

- Optimal control: problem formulation
- Discrete Mechanics and Optimal Control
- Preservation properties and approximation error

briefly:

- Variational approach to multirate integration
- related works:

Betsch, Bock, Bonnans, Diehl, Gerdts, Hager, Kobilarov, Leok, Leyendecker, Marsden, Ortiz, von Stryk, West, Chyba, Hairer, Vilmart, ...



Optimal Control Problem

$$\min_{x(\cdot),u(\cdot),(T)} J(x,u) = \int_0^T C(x(t),u(t)) dt + \Phi(x(T))$$
subject to



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Necessary optimality conditions: Pontryagin maximum principle

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Necessary optimality conditions

Theorem (Pontryagin Maximum Principle)

Let (x^*, u^*) be an optimal solution. Then, there exists a function $\lambda : [0, T] \to \mathbb{R}^{n_x}$ and a vector $\alpha \in \mathbb{R}^{n_r}$ such that

$$\begin{aligned} \mathcal{H}(x^*(t), u^*(t), \lambda(t)) &= \max_{u(t) \in U} \mathcal{H}(x(t), u(t), \lambda(t)) \,\forall t \in [0, T], \\ \dot{x}^*(t) &= \nabla_\lambda \mathcal{H}(x^*(t), u^*(t), \lambda(t)), \, x^*(0) = x_0, \\ \dot{\lambda}(t) &= -\nabla_x \mathcal{H}(x^*(t), u^*(t), \lambda(t)), \\ \lambda(T) &= \nabla_x \Phi(x^*(T)) - \nabla_x r(x^*(T)) \,\alpha \end{aligned}$$

with the Hamiltonian $\mathcal{H}(x(t), u(t), \lambda(t)) = -C(x(t), u(t)) + \lambda^{T}(t) \cdot f(x(t), u(t)).$



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with the Hamiltonian $\mathcal{H}(x(t), u(t), \lambda(t)) = -C(x(t), u(t)) + \lambda^{T}(t) \cdot f(x(t), u(t)).$



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with the Hamiltonian $\mathcal{H}(x(t), u(t), \lambda(t)) = -C(x(t), u(t)) + \lambda^{T}(t) \cdot f(x(t), u(t)).$



Mechanical system: variational formulation

- n-dimensional configuration manifold Q
- ▶ Lagrangian $L: TQ \to \mathbb{R}, \ L(q, \dot{q}) = K(q, \dot{q}) V(q)$



Lagrange-d'Alembert principle

$$\delta \int_0^T L(q(t), \dot{q}(t)) \, \mathrm{d}t + \int_0^T f_L(q(t), \dot{q}(t), u(t)) \cdot \delta q(t) \, \mathrm{d}t = 0$$

 $\Rightarrow \text{Euler-Lagrange equations} \\ \frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial}{\partial \dot{q}} L(q(t), \dot{q}(t)) - \frac{\partial}{\partial q} L(q(t), \dot{q}(t)) = f_L(q(t), \dot{q}(t), u(t))$



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 $\Rightarrow \text{Euler-Lagrange equations} \\ \frac{d}{dt} \frac{\partial}{\partial \dot{q}} L(q(t), \dot{q}(t)) - \frac{\partial}{\partial q} L(q(t), \dot{q}(t)) = f_L(q(t), \dot{q}(t), u(t))$

 $\Rightarrow \dot{x}(t) = f(x(t), u(t))$ with $x = (q, \dot{q})$



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What are the relevant preservation properties?

Noether Theorem (no forcing)

Invariance of the Lagrangian under a group action of a Lie group G leads to the **preservation of momentum maps** along the trajectory.

- ▶ Lagrangian $L: TQ \rightarrow \mathbb{R}$, Lie group G, left action $\Psi_g: Q \rightarrow Q, g \in G$
- invariance of the Lagrangian: $L \circ \Psi_g^{TQ} = L$
- ▶ left action induces momentum map J with $J \circ F_L^t = J$ for all t with Lagrangian flow $F_L^t : TQ \to TQ$



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Symplecticity (no forcing)

Lagrangian flow is symplectic: $(F_L^t)^*(\Omega) = \Omega$.



Replace

- state space TQ with $Q \times Q$
- curve $q(t) \in Q$ with sequence $q_d = \{q_k\}_{k=0}^N$
- curve $u(t) \in U$ with sequence $u_d = \{u_k\}_{k=0}^{N-1}$





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Approximate Lagrangian and virtual work:

• discrete Lagrangian $L_d: Q imes Q
ightarrow \mathbb{R}$

$$L_d(q_k,q_{k+1}) pprox \int_{kh}^{(k+1)h} L(q(t),\dot{q}(t)) \mathrm{d}t$$

- ► Discrete Lagrange-d'Alembert principle $\delta \sum_{k=0}^{N-1} L_d(q_k, q_{k+1}) + \sum_{k=0}^{N-1} \left[f_k^- \cdot \delta q_k + f_k^+ \cdot \delta q_{k+1} \right] = 0$
- Discrete Euler-Lagrange equations (DEL)

 $D_1L_d(q_k, q_{k+1}) + D_2L_d(q_{k-1}, q_k) + f_k^- + f_{k-1}^+ = 0$

- Variational integrator for forward simulation
 - ⊕ symplectic (preservation of symplectic form) ⇒ good
 long-time energy behavior
 HAIRER & LUBICH (2004)
 - ⊕ preserves momentum maps, if L_d is invariant under group action
 MARSDEN & WEST (2001)
 - order given by order of discrete Lagrangian and forces
- boundary conditions: continuous boundary conditions on TQ are transformed in discrete boundary conditions on Q × Q via the discrete Legendre transformation



Discrete optimal control problem

$$\begin{array}{l} (\text{DMOC: OB, JUNGE \& MARSDEN (2008)})\\ \min_{q_d, u_d, (h)} J_d(q_d, u_d) = \sum_{k=0}^{N-1} C_d(q_k, q_{k+1}, u_k) + \Phi_d(q_{N-1}, q_N, u_{N-1})\\ \text{s.t. DEL} \quad D_2 L_d(q_{k-1}, q_k) + D_1 L_d(q_k, q_{k+1}) + f_{k-1}^+ + f_k^- = 0\\ \text{initial conditions} \qquad s_d(q_0, q_1, u_0, q^0, \dot{q}^0) = 0\\ \text{path constraints} \qquad h_d(q_k, q_{k+1}, u_k) \geq 0\\ \text{final constraint} \qquad r_d(q_{N-1}, q_N, u_{N-1}, q^T, \dot{q}^T) = 0 \end{array}$$



Discrete optimal control problem

$$(DMOC: OB, JUNGE \& MARSDEN (2008))$$
$$\min_{q_d, u_d, (h)} J_d(q_d, u_d) = \sum_{k=0}^{N-1} C_d(q_k, q_{k+1}, u_k) + \Phi_d(q_{N-1}, q_N, u_{N-1})$$
s.t. DEL $D_2L_d(q_{k-1}, q_k) + D_1L_d(q_k, q_{k+1}) + f_{k-1}^+ + f_k^- = 0$ initial conditions $s_d(q_0, q_1, u_0, q^0, \dot{q}^0) = 0$ path constraints $h_d(q_k, q_{k+1}, u_k) \ge 0$ final constraint $r_d(q_{N-1}, q_N, u_{N-1}, q^T, \dot{q}^T) = 0$
$$\bigcup$$
restricted optimization problem $[\min_{\xi} \varphi(\xi) \text{ s. t. } a(\xi) = 0, \ b(\xi) \ge 0]$



Discrete optimal control problem

(DMOC: OB, JUNGE & MARSDEN (2008))

$$\min_{q_d, u_d, (h)} J_d(q_d, u_d) = \sum_{k=0}^{N-1} C_d(q_k, q_{k+1}, u_k) + \Phi_d(q_{N-1}, q_N, u_{N-1})$$
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 ψ
restricted optimization problem
 $\boxed{\min_{\xi} \varphi(\xi) \text{ s. t. } a(\xi) = 0, b(\xi) \ge 0}$
 ψ
necessary optimality conditions
Karush-Kuhn-Tucker: $\nabla \varphi(\xi) - \lambda^T \nabla a(\xi) - \mu^T \nabla b(\xi) = 0$
with Lagrange multipliers λ und μ .

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Structure preservation in presence of forcing

Consequence from Noether:

- evolution of momentum map exactly determined by control forces
- ► preservation of momentum map if control forces act orthogonal to group action $(\langle f_L, \xi_Q \rangle = 0)$



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Example: control force acts orthogonal to relative velocity \Rightarrow exact preservation of velocity's absolute value





The adjoint system





The adjoint system



 RK methods: order of adjoint scheme is in general NOT the same as for state scheme [HAGER 2000], [BONNANS, LAURENT-VARIN 2006]



The adjoint system



- RK methods: order of adjoint scheme is in general NOT the same as for state scheme
 - [Hager 2000], [Bonnans, Laurent-Varin 2006]
- BUT: using the *right* scheme for the state system, state and adjoint scheme coincide (e.g. special class of VI)



Properties of DMOC: order of approximation

Proof strategy:

1 VI of order r given by $L_d = h \sum_{i=1}^{s} b_i L(Q_i, \dot{Q}_i)$

symplectic partitioned Runge-Kutta scheme of order r

≏

[Suris 1990], [Marsden, West 2001]



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[Suris 1990], [Marsden, West 2001]

- 2 determine approximation order κ of adjoint system (as for RK methods [HAGER 2000]):
 - necessary optimality conditions for continuous system (Pontryagin)
 - necessary optimality conditions for discrete system (KKT)



1 Hamiltonian formulation

- Lagrangian (q, \dot{q}) can be rewritten as Hamiltonian (q, p)
- ▶ partitioned Runge-Kutta scheme [MARSDEN, WEST 2001]

$$\begin{aligned} q_1 &= q_0 + h \sum_{j=1}^s b_j^q \dot{Q}_j, \qquad p_1 = p_0 + h \sum_{j=1}^s b_j^p \dot{P}_j, \\ Q_i &= q_0 + h \sum_{j=1}^s a_{ij}^q \dot{Q}_j, \qquad P_i = p_0 + h \sum_{j=1}^s a_{ij}^p \dot{P}_j \\ P_i &= \frac{\partial L}{\partial \dot{q}} (Q_i, \dot{Q}_i), \qquad \dot{P}_i = \frac{\partial L}{\partial q} (Q_i, \dot{Q}_i) + f_L (Q_i, \dot{Q}_i, U_i) \end{aligned}$$

i = 1, ..., s, with internal stages (Q_i, P_i) , control samples $U_i = u_d(t_0 + c_i h)$, symplectic if $b_i^q a_{ij}^p + b_j^p a_{ji}^q = b_i^q b_j^p$, $b_i^q = b_i^p$.



2 Continuous adjoint system

- substitute optimal control u(q(t), p(t), λ^q(t), λ^p(t)) obtained by solving minimization problem
- two-point boundary problem

$$\begin{split} \dot{q}(t) &= \nu(q(t), p(t)), \quad q(0) = q^{0}, \\ \dot{p}(t) &= \eta(q(t), p(t), \lambda^{q}(t), \lambda^{p}(t)), \quad p(0) = p^{0}, \\ \dot{\lambda}^{q}(t) &= \phi^{q}(q(t), p(t), \lambda^{q}(t), \lambda^{p}(t)), \quad \lambda^{q}(T) = \Psi^{q}(q(T), p(T)), \\ \dot{\lambda}^{p}(t) &= \phi^{p}(q(t), p(t), \lambda^{q}(t), \lambda^{p}(t)), \quad \lambda^{p}(T) = \Psi^{p}(q(T), p(T)). \end{split}$$



2 Discrete adjoint system (I)

$$\begin{aligned} q_{k+1} &= q_k + h \sum_{i=1}^{s} b_i \nu(Q_{ki}, P_{ki}), \quad q_0 = q^0, \\ p_{k+1} &= p_k + h \sum_{i=1}^{s} b_i \eta(Q_{ki}, P_{ki}, \chi_{ki}^q, \chi_{ki}^p), \quad p_0 = p^0, \\ \lambda_{k+1}^q &= \lambda_k^q + h \sum_{i=1}^{s} b_i \phi^q(Q_{ki}, P_{ki}, \chi_{ki}^q, \chi_{ki}^p), \quad \lambda_N^q = \Psi^q(q_N, p_N), \\ \lambda_{k+1}^p &= \lambda_k^p + h \sum_{i=1}^{s} b_i \phi^p(Q_{ki}, P_{ki}, \chi_{ki}^q, \chi_{ki}^p), \quad \lambda_N^p = \Psi^p(q_N, p_N), \end{aligned}$$



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2 Discrete adjoint system (II)

$$\begin{aligned} Q_{ki} &= q_k + h \sum_{j=1}^{s} a_{ij}^{q} \nu(Q_{kj}, P_{kj}), \\ P_{ki} &= p_k + h \sum_{j=1}^{s} a_{ij}^{p} \eta(Q_{kj}, P_{kj}, \chi_{kj}^{q}, \chi_{kj}^{p}), \\ \chi_{ki}^{q} &= \lambda_k^{q} + h \sum_{j=1}^{s} \bar{a}_{ij}^{q} \phi^{q}(Q_{kj}, P_{kj}, \chi_{kj}^{q}, \chi_{kj}^{p}), \\ \chi_{ki}^{p} &= \lambda_k^{p} + h \sum_{j=1}^{s} \bar{a}_{ij}^{p} \phi^{p}(Q_{kj}, P_{kj}, \chi_{kj}^{q}, \chi_{kj}^{p}), \\ \bar{a}_{ij}^{q} &= \frac{b_i b_j - b_j a_{ji}^{q}}{b_i}, \quad \bar{a}_{ij}^{p} = \frac{b_i b_j - b_j a_{ji}^{p}}{b_i}. \end{aligned}$$

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2 Order of approximation

Theorem (Order of approximation)

If the partitioned symplectic Runge-Kutta discretization of the state system is of order κ and $b_i > 0$ for each i, then the scheme for the adjoint system is again a partitioned symplectic Runge-Kutta scheme of the same order (in particular we obtain the same schemes for (q, p) and (λ^p, λ^q)).

Proof: With the symplecticity condition $a_{ij}^p = \frac{b_i b_j - b_j a_{ji}^q}{b_i}$ it holds

$$ar{a}^q_{ij} = a^p_{ij} \ ar{a}^p_{ij} = a^q_{ij}$$

[OB, JUNGE, MARSDEN 2009]



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Properties of DMOC: convergence rates



- two-link pendulum optimally (minimal control effort) controlled from lower to upper equilibrium
- same convergence rates for configuration, control and adjoint

DMOC and ADOI (-log(max(u-u_)) rates control convergence 2 -log(h) - DMOC DMOC and ADOL $-\log(\max(\lambda - \lambda))$ adjoint convergence rates -log(h)

DMOC

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...to give answers to ?

A+B optimal control problem + structure preserving integration

- inherited from integration scheme (var. formulation): evolution of momentum maps due to symmetries determined by control forces
- gain for optimal control problems (symplecticity): order of adjoint system = order of state system



...to give answers to ?

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extension to (holonomic) constrained systems



Optimal control of multi-body dynamics (DMOCC) [LEYENDECKER, OB, MARSDEN, ORTIZ 08]

- ▶ constraint manifold $\mathcal{C} = \{q \, | \, q \in Q, g(q) = 0\} \subset Q$
- ► discrete version of Lagrange-d'Alembert principle with augmented Lagrangian \$\overline{L}(q, \overline{q}, \mu) = L(q, \overline{q}) - g^T(q) \cdot \mu\$
- Forced constrained discrete EL equations

$$\begin{bmatrix} D_2 L_d(q_{k-1}, q_k) + D_1 L_d(q_k, q_{k+1}) - h G^T(q_k) \cdot \mu_k + f_{k-1}^+ + f_k^- \end{bmatrix} = 0$$

$$g(q_{k+1}) = 0$$



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- Forced constrained discrete EL equations

$$P^{T}(q_{k}) \left[D_{2}L_{d}(q_{k-1}, q_{k}) + D_{1}L_{d}(q_{k}, q_{k+1}) - hG^{T}(q_{k}) \cdot \mu_{k} + f_{k-1}^{+} + f_{k}^{-} \right] = 0$$

$$g(q_{k+1}) = 0$$

- dimension reduction via discrete nullspace matrix, relative reparametrization of configurations and forces
- \oplus reduced scheme with minimal dimension
- $_\oplus$ symplectic and momentum consistent
- $_{\oplus}$ exact constraint fulfillment



Example (DMOCC): satellite reorientation

[Leyendecker, OB, Marsden, Ortiz 08]

- satellite actuated via rotors attached at main body
- Goal: attitude control with minimal control effort
- preservation of total angular momentum





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Variational multirate integration

(joint work with S. LEYENDECKER)

- mechanical systems with dynamics on different time scales
- Lagrangian

$$egin{aligned} \mathcal{L}(q,\dot{q}) &= \mathcal{T}(\dot{q}) - \mathcal{U}(q), \quad \mathcal{U}(q) &= \mathcal{V}(q) + \mathcal{W}(q) \end{aligned}$$
 with $\mathcal{W}(q) &= rac{1}{arepsilon}ar{\mathcal{W}}(q), \ arepsilon \ll 1. \end{aligned}$

- ▶ assumption: separation in slow and fast variables $q = (q^s, q^f)$
- ▶ fast potential W only dependent on fast variables q^f
- Lagrangian

$$L(q,\dot{q})=T(\dot{q})-V(q)-W(q^{f})$$

 aim: resolve slow and fast dynamics but with less function evaluations and smaller set of variables



Multirate: discrete approximation

consider two time grids



- discrete slow variables $\{q_k^s\}_{k=0}^N$ with $q_k^s \approx q^s(t_k)$
- discrete fast variables $\{\{q_k^{f,m}\}_{m=0}^p\}_{k=0}^{N-1}$ with $q_k^{f,p} = q_{k+1}^{f,0}$ and $q_k^{f,m} \approx q^f(t_k^m)$



Discrete Lagrangian

$$\begin{split} \bar{L}_{d}(q_{k}^{s}, q_{k+1}^{s}, q_{k}^{f,0}, \dots, q_{k}^{f,p}) &= \\ \sum_{m=0}^{p-1} L_{d}(q_{k}^{s}, q_{k+1}^{s}, q_{k}^{f,m}, q_{k}^{f,m+1}) \\ &= \sum_{m=0}^{p-1} \left[T_{d}(q_{k}^{s}, q_{k+1}^{s}, q_{k}^{f,m}, q_{k}^{f,m+1}) - V_{d}(q_{k}^{s}, q_{k+1}^{s}, q_{k}^{f,m+1}) - V_{d}(q_{k}^{s}, q_{k+1}^{s}, q_{k}^{f,m+1}) - W_{d}(q_{k}^{f,m}, q_{k}^{f,m+1}) \right] \\ \end{split}$$
with $L_{d} = T_{d} - V_{d} - W_{d}$

$$\begin{split} & \underset{t_{k-1}}{\overset{t_{k-1}}{\longrightarrow} \overset{t_{m-1}}{\longrightarrow} \overset{t_{m-1$$



discrete action

$$\mathfrak{S}_{d}\left(\{q_{k}^{s}\}_{k=0}^{N},\{\{q_{k}^{f,m}\}_{m=0}^{p}\}_{k=0}^{N-1}\right) = \sum_{k=0}^{N-1} \sum_{m=0}^{p-1} L_{d}(q_{k}^{s},q_{k+1}^{s},q_{k}^{f,m},q_{k}^{f,m+1})$$

stationary discrete action

$$\delta\mathfrak{S}_{d} = \delta \sum_{k=0}^{N-1} \sum_{m=0}^{p-1} L_{d}(q_{k}^{s}, q_{k+1}^{s}, q_{k}^{f,m}, q_{k}^{f,m+1})$$

$$= \sum_{k=0}^{N-1} \sum_{m=0}^{p-1} D_{1}L_{d} \cdot \delta q_{k}^{s} + D_{2}L_{d} \cdot \delta q_{k+1}^{s}$$

$$+ \sum_{k=0}^{N-1} \sum_{m=0}^{p-1} D_{3}L_{d} \cdot \delta q_{k}^{f,m} + D_{4}L_{d} \cdot \delta q_{k}^{f,m+1} = 0$$

gives discrete Euler-Lagrange equations for fast and slow dynamics $% \left({{{\left[{{{\left[{{{c_{{\rm{m}}}}} \right]}} \right]}_{\rm{max}}}}} \right)$

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Triple pendulum in 3D

- different quadrature rules lead to different schemes
- invariance of discrete
 Lagrangian: preserved
 momentum maps
- discrete sympletic form preserved





Conclusion and future directions

Conclusion

 fully discrete variational formulation for the optimal control of mechanical systems

variational: structure preserving symplectic: approximation order of adjoint system = approximation order of state system → order of adjoint system also for extended DMOC versions (with holonomic constraints)

variational multirate integration

- unified variational approach for the simulation of mechanical systems with different time scales
- structure preserving

 \rightarrow analysis: comparison to existing methods, long-time behavior, efficiency \rightarrow use for optimal control

