

Numerical solution of the wave equation with acoustic boundary conditions

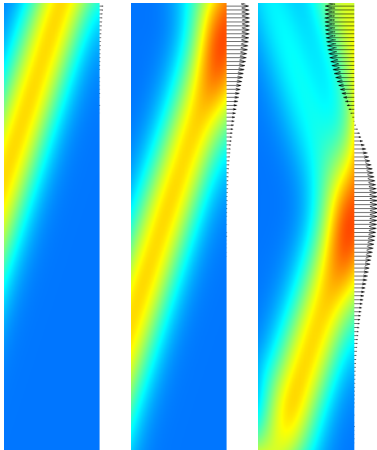
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Numerical Analysis Group



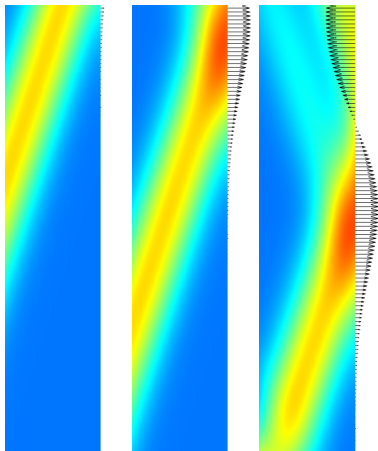
CRC 1173

Wave
phenomena



example

- design of concert hall
- model walls as tiny springs
- springs react to excess pressure



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setting

- bounded domain $\Omega \subset \mathbb{R}^2$
- boundary $\Gamma = \partial\Omega$
- u velocity potential
- δ infinitesimal normal displ. of Γ
- Γ not moving!

Acoustic boundary conditions

introduced in [Beale, Rosencrans '74]

- $u: [0, T] \times \Omega \rightarrow \mathbb{R}$
- $\delta: [0, T] \times \Gamma \rightarrow \mathbb{R}$

wave equation with acoustic b.c.

u and δ are governed by:

$$\partial_t^2 u = c^2 \Delta u \quad \text{in } \Omega$$

$$m \partial_t^2 \delta + k \delta + d \delta = -\rho \partial_t u \quad \text{on } \Gamma$$

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+ initial conditions for $u, \partial_t u, \delta, \partial_t \delta$

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- dynamic boundary condition

Boundary conditions for the wave equation

■ energies of u

kinetic $KE(u) = \frac{1}{2} \int_{\Omega} \frac{1}{c^2} |\partial_t u|^2 dx$

potential $PE(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx$

[G.R. Goldstein '06] for details

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minimizing action functional $S(u) = \int_0^T KE(u) - PE(u) dt$

$$\partial_t^2 u = c^2 \Delta u \quad \text{in } \Omega$$

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- kinetic boundary conditions
- Robin and Neumann b.c. are contained
- dynamic boundary condition for $\lambda > 0$

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Comparison of different boundary conditions

Neumann	Robin	kinetic	acoustic
$\partial_\nu u = 0$	$\partial_\nu u + u = 0$	$\partial_t^2 u + ku = -\partial_\nu u$	$\partial_t^2 \delta + k\delta = -\partial_t u$

Analysis I

well-posedness of wave equation

$$\partial_t^2 u = \Delta u \quad \text{in } \Omega \quad + \quad \text{i.c.}$$

with Neumann b.c.

$$\partial_\nu u = 0 \quad \text{on } \Gamma$$

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$$\partial_t^2 \delta + k\delta = -\partial_t u \quad \text{on } \Gamma$$

as 1st order evolution equation

$$\frac{d}{dt} \vec{u}(t) = A\vec{u}(t) \quad + \quad \text{i.c.}$$

with operator

$$A\vec{u} = \begin{bmatrix} 0 & I \\ \Delta & 0 \end{bmatrix} \begin{bmatrix} u \\ u_t \end{bmatrix} \quad \mathcal{A}\vec{u} = \begin{bmatrix} u_t \\ \Delta u \\ \delta_t \\ -\gamma(u_t) - k\delta \end{bmatrix} \quad \vec{u} = \begin{bmatrix} u \\ u_t \\ \delta \\ \delta_t \end{bmatrix}$$

Analysis II

- (energy) Hilbert space

$$H = H_c^1(\Omega) \times L^2(\Omega) \times L^2(\Gamma) \times L^2(\Gamma) = \mathcal{H}$$

$$\|\vec{u}\|_H^2 = \frac{1}{2} \int_{\Omega} |\nabla u|^2 + |u_t|^2 \, dx + \frac{1}{2} \int_{\Gamma} k |\delta|^2 + |\delta_t|^2 \, d\sigma = \|\vec{u}\|_{\mathcal{H}}^2$$

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- domain of operator

$$\mathcal{D}(\mathcal{A}) = \{\vec{u} \in H \mid \Delta u \in L^2(\Omega), u_t \in H^1(\Omega), \partial_\nu u = 0 \text{ on } \Gamma\}$$

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Theorem ([Beale '76, Thm. 2.1])

- \mathcal{A} closed, densely defined and skewadjoint in \mathcal{H}



$$\frac{d}{dt} \vec{u}(t) = \mathcal{A} \vec{u}(t), \quad \vec{u}(0) = \vec{u}_0 \in \mathcal{D}(\mathcal{A})$$

is well-posed and $\|\vec{u}\|_{\mathcal{H}} = \text{const.}$

method of lines

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2. numerical time integration of stiff ODE

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- **solution:** choose $\mathcal{H} = H^1(\Omega) \times L^2(\Omega) \times L^2(\Gamma) \times L^2(\Gamma)$

$$\|\vec{u}\|_{\mathcal{H}}^2 = \int_{\Omega} |u|^2 \, dx + \|\vec{u}\|_{\mathcal{H}_c}^2$$

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- $\mathcal{A}: \mathcal{H} \supset \mathcal{D}(\mathcal{A}) \rightarrow \mathcal{H}$ not skewadjoint, but

$$\left(\vec{v} \left| \left(\mathcal{A} - \frac{1}{2} \text{I} \right) \vec{v} \right. \right)_{\mathcal{H}} \leq 0, \quad \vec{v} \in \mathcal{D}(\mathcal{A})$$

Corollary

$\mathcal{A}: \mathcal{D}(\mathcal{A}) \rightarrow \mathcal{H}$ is the infinitesimal generator of a C_0 -semigroup

Variational formulation

Green's formula: $\Delta u \in L^2(\Omega)$, $v_t \in H^1(\Omega)$, $\partial_\nu u = \delta_t$

$$\begin{aligned}
 (\vec{v} | \mathcal{A}\vec{u})_{\mathcal{H}} &= (v | u_t)_{1,\Omega} + (v_t | \Delta u)_{0,\Omega} \\
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variational problem

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Spatial discretization

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semidiscrete problem

find $\vec{u}_h: [0, T] \rightarrow \mathcal{V}_h$ s.t.

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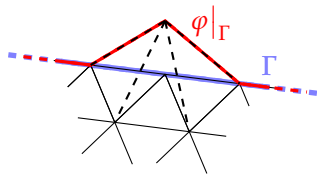
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finite element spaces

- \mathcal{T}_h regular triangulation of Ω
- $V_h^\Omega =$ pcw. linear FEs over \mathcal{T}_h
- $V_h^\Gamma = \gamma(V_h^\Omega)$



Convergence FE-semidiscretization I

1. split error with orthogonal projection $\mathcal{P}_h: \mathcal{H} \rightarrow \mathcal{V}_h$

$$\vec{e} = \vec{u}_h - \vec{u} = (\vec{u}_h - \mathcal{P}_h \vec{u}) + (\mathcal{P}_h \vec{u} - \vec{u}) = \vec{e}_h + \vec{e}_p$$

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3. set $\vec{v}_h = \vec{e}_h$

$$\frac{1}{2} \frac{d}{dt} \|\vec{e}_h\|_{\mathcal{H}}^2 = s(\vec{e}_h, \vec{e}_h) + s(\vec{e}_h, \vec{e}_p) \leq C \|\vec{e}_h\|_{\mathcal{H}}^2 + s(\vec{e}_h, \vec{e}_p)$$

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4. with \mathcal{P}_h -properties and standard FE approximation results

$$\begin{aligned} |s(\vec{e}_h, \vec{e}_p)| &\leq C \|\vec{e}_h\|_{\mathcal{H}}^2 + C \left(\|P_{0,\Omega} u_t - u_t\|_{1,\Omega}^2 + \|P_{1,\Omega} u - u\|_{0,\Omega}^2 \right) \\ &\leq C \|\vec{e}_h\|_{\mathcal{H}}^2 + \tilde{C} (|u_t|_{2,\Omega}, |u|_{2,\Omega}) h^2 \end{aligned}$$

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$$\frac{1}{2} \frac{d}{dt} \|\vec{e}_h\|_{\mathcal{H}}^2 \leq C \|\vec{e}_h\|_{\mathcal{H}}^2 + \tilde{C} h^2$$

6. apply Gronwall's lemma

$$\|\vec{e}_h(t)\|_{\mathcal{H}} \leq Ch$$

4. with \mathcal{P}_h -properties and standard FE approximation results

$$\begin{aligned} |s(\vec{e}_h, \vec{e}_P)| &\leq C \|\vec{e}_h\|_{\mathcal{H}}^2 + C \left(\|P_{0,\Omega} u_t - u_t\|_{1,\Omega}^2 + \|P_{1,\Omega} u - u\|_{0,\Omega}^2 \right) \\ &\leq C \|\vec{e}_h\|_{\mathcal{H}}^2 + \tilde{C} (|u_t|_{2,\Omega}, |u|_{2,\Omega}) h^2 \end{aligned}$$

C and \tilde{C} independent of h

5. then

$$\frac{1}{2} \frac{d}{dt} \|\vec{e}_h\|_{\mathcal{H}}^2 \leq C \|\vec{e}_h\|_{\mathcal{H}}^2 + \tilde{C} h^2$$

6. apply Gronwall's lemma

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7. consider full error

$$\|\vec{e}(t)\|_{\mathcal{H}} \leq \|\vec{e}_h(t)\|_{\mathcal{H}} + \|\vec{e}_P(t)\|_{\mathcal{H}} \leq Ch$$

Theorem (H., Hochbruck 2015)

$$\vec{u} \in L^\infty\left([0, T]; H^2(\Omega) \times H^1(\Omega) \times H^1(\Gamma) \times H^1(\Gamma)\right)$$

$$u_t \in L^2\left([0, T]; H^2(\Omega)\right)$$

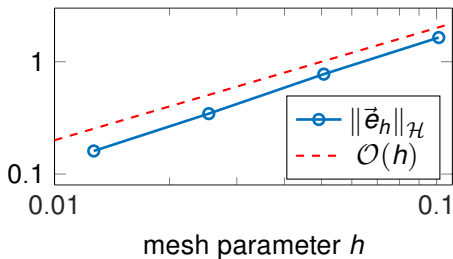
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Comment on time integration

method of lines gives stiff problem

$$\frac{d}{dt} \mathbf{M}_h \mathbf{u}(t) = \mathbf{S}_h \mathbf{u}(t)$$

with exact solution

$$\mathbf{u}(t) = \exp\left(t \mathbf{M}_h^{-1} \mathbf{S}_h\right) \mathbf{u}(0), \quad t \geq 0$$

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 - explicit time stepping schemes
 - direct approximation in Krylov subspaces
2. rational approximation
 - implicit time stepping schemes
 - direct approximation with rational Krylov methods

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1. polynomial approximation
 - explicit time stepping schemes
 - direct approximation in Krylov subspaces
2. rational approximation
 - implicit time stepping schemes
 - direct approximation with rational Krylov methods
3. combinations of both

- evolution equation as boundary conditions, e.g.

$$\begin{aligned}\partial_t^2 u &= c^2 \Delta u && \text{in } \Omega \\ \partial_t^2 \delta + c_\Gamma \Delta_\Gamma \delta &= -\partial_t u && \text{on } \Gamma \\ \partial_t \delta &= \partial_\nu u && \text{on } \Gamma\end{aligned}$$

- non-linear boundary conditions and coupling
- splitting methods for domain and boundary part

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